

Blasius Theorem :

Statement : In steady two dimensional motion given by the complex potential $w = f(z) = \phi + i\psi$, if the pressure thrusts on the fixed cylinder of any shape are represented by a force (X, Y) and a couple of moment M about the origin of co-ordinates, then neglecting external forces and

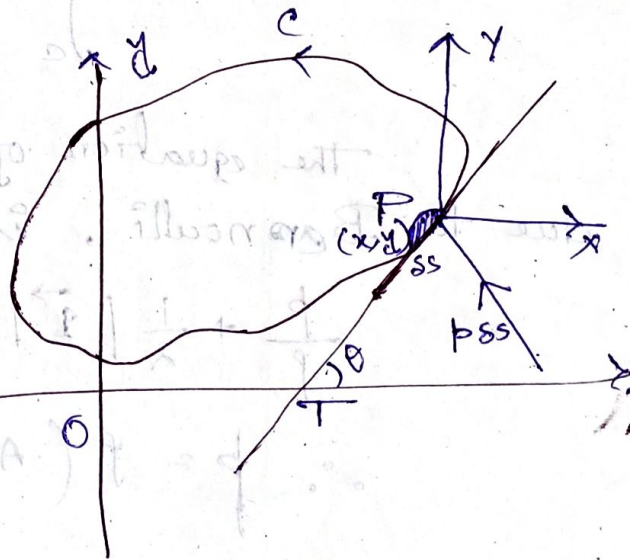
$$X - iY = \frac{i\rho}{2} \int_C \left(\frac{dw}{dz} \right)^2 dz \quad \text{and}$$

$$M = \text{Re} \left[-\frac{\rho}{2} \int_C \left(\frac{dw}{dz} \right)^2 z dz \right]$$

where ρ is the density and integrals are taken round the contour C of the cylinder.

Proof :

Let us consider an element δs about the point $P(x, y)$ of the fixed cylinder.



Let C be the boundary of the cylinder of any shape and size

Let the tangent PT makes an angle θ with x axis.
The inward normal at P makes an angle $\left(\frac{\pi}{2} + \theta \right)$ with x axis

Then the thrust components at P are
 $p s s \cos(\alpha_2 + \theta)$ along x axis and
 $p s s \sin(\alpha_2 + \theta)$ along y axis
 i.e., $-p s s \sin \theta$ along x axis and
 $p s s \cos \theta$ along y axis.

$$\therefore X = \sum -p s s \sin \theta = - \int_C p \sin \theta ds$$

$$\text{and } Y = \sum p s s \cos \theta = \int_C p \cos \theta ds$$

$$\begin{aligned} \therefore X - iY &= \int_C p (-\sin \theta - i \cos \theta) ds \\ &= -i \int_C p (\cos \theta - i \sin \theta) ds \end{aligned}$$

The equation of steady state mol
 due to Bernoulli, is

$$\frac{p}{\rho} + \frac{1}{2} |\vec{q}|^2 = \text{Constant} = A,$$

$$\therefore p = \rho \left(A - \frac{1}{2} |\vec{q}|^2 \right)$$

$$\text{Hence } X - iY = -i\rho \int_C \left(A - \frac{1}{2} |\vec{q}|^2 \right) ds$$

$$= \frac{1}{2} i\rho \int_C |\vec{q}|^2 e^{-i\theta} ds$$

$$-i\rho A \int_C e^{-i\theta} ds$$

$$2^{\text{nd}} \text{ integral} = -i\phi A \int_c (\cos\theta ds - i\sin\theta ds)$$

$$= -i\phi A \int_c (dx - idy)$$

$$= -i\phi A \int_c d\bar{z}$$

= 0, by Cauchy's theorem, as c is a closed contour.

$$\tan\theta = \frac{dy}{dx}$$

$$\sin\theta = \frac{dy}{ds}$$

$$\cos\theta = \frac{dx}{ds}$$

$$z = x + iy$$

Hence,

$$x - iy = \frac{1}{2} i\phi \int_c |\vec{q}|^2 e^{-i\theta} ds$$

Let u, v be the components of the vector \vec{q} .
 $|\vec{q}|^2 = u^2 + v^2$

and we know that

$$\frac{dw}{dz} = -u + iv$$

$$= -|\vec{q}| \cos\theta + i|\vec{q}| \sin\theta$$

$$u = -\frac{\partial\phi}{\partial x}$$

$$v = \frac{\partial\phi}{\partial y}$$

$$\text{or } \vec{q} = -\nabla\phi$$

and c.

$$= -|\vec{q}| (\cos\theta - i\sin\theta)$$

$$= -|\vec{q}| e^{-i\theta}$$

$$\left(\frac{dw}{dz}\right)^2 = |\vec{q}|^2 e^{-2i\theta}$$

$$x - iy = \left(\frac{dw}{dz}\right)^2 dz$$

$$= |\vec{q}|^2 e^{-2i\theta} (dx + idy)$$

$$= |\vec{q}|^2 e^{-2i\theta} (\cos\theta ds + i\sin\theta ds)$$

$$= |\vec{q}|^2 e^{-2i\theta} e^{i\theta} ds$$

$$= |\vec{q}|^2 e^{-i\theta} ds$$

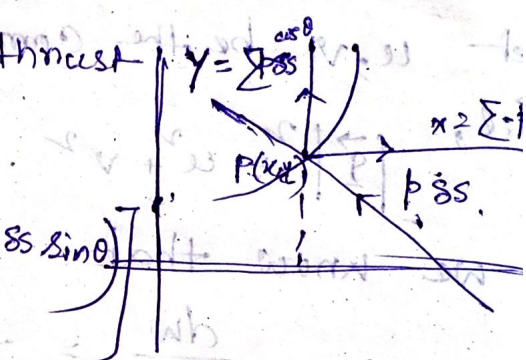
$$\therefore x - iy = \frac{1}{2} i p \int_C |\vec{q}|^2 e^{-i\theta} ds$$

$$= \frac{ip}{2} \int_C \left(\frac{dw}{dz} \right)^2 dz$$

which is the first result and moment of the thrust about the origin

Now the moment of thrust about origin O is

$$M = \int_C \left[(+p s \cos\theta) x - y (-p s \sin\theta) \right]$$



$$= \int_C (x p \cos\theta + y p \sin\theta) ds$$

$$= \int_C p \left(A - \frac{1}{2} |\vec{q}|^2 \right) (x dx + y dy)$$

$$= p A \int_C \frac{d(x^2 + y^2)}{2} - \frac{p}{2} \int_C |\vec{q}|^2 (x dx + y dy)$$

Now, $\int_C d(x^2 + y^2) = 0$ (by Cauchy theorem).

Hence

$$N = -\frac{1}{2} \int_C |\vec{q}|^2 (x dx + y dy)$$

$$= \operatorname{Re} \left[\frac{1}{2} \int_C |\vec{q}|^2 z e^{-i\theta} ds \right]$$

$$\begin{aligned} \therefore z e^{-i\theta} ds &= (x + iy) (\cos \theta ds - i \sin \theta ds) \\ &= (x + iy) (dx - i dy) \\ &= (x dx + y dy) + i (y dx - x dy) \\ \Rightarrow x dx + y dy &= \operatorname{Re} \left\{ z e^{-i\theta} ds \right\} \end{aligned}$$

$$\therefore N = \operatorname{Re} \left[-\frac{1}{2} \int_C z \left(\frac{dw}{dz} \right)^2 dz \right]$$

Hence the theorem.

Ex: A source of fluid situated in space of two dimension is of such strength that

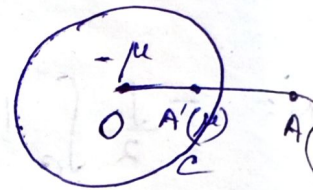
$2\pi\mu$ represents that the mass of fluid of density ρ emitted for unit of time

Show that the force necessary to hold a circular disc at rest in the plane source is

$\frac{2\pi\mu\rho a^2}{r(r^2 - a^2)}$, where a is the radius of the disc and r be the distance of the source from its centre. In what

direction is the disc urged by the pressure.

Solⁿ: Let x and y be the components of the required force. Then by Blasius theorem,



$$X - iy = \frac{i\rho}{2} \int_C \left(\frac{dw}{dz} \right)^2 dz$$

, where C represents the boundary of the disc.

Since $2\rho\mu$ represents the mom of the fluid emitted,

hence strength of the source is μ

Thus the image of a source ($+\mu$) at A where $(OA = r)$, is a source $+\mu$ at the inverse point at A' s.t.

$$OA \cdot OA' = a^2$$

and a sink ($-\mu$) at O .

Now, $OA \cdot OA' = a^2$ gives $OA' = \frac{a^2}{r} = r'$, say

\therefore The complex potential due to object system with rigid boundary is equivalent to the complex potential due to the object system and its image system with no rigid boundary.

Hence $W = -\mu \log(2-n) - \mu \log(2-n') + \mu$

$$\therefore \frac{dw}{dz} = -\frac{\mu}{2-n} - \frac{\mu}{2-n'} + \left(\frac{\mu}{2} + \frac{1}{n'}\right)$$

$$= -\mu \left[\frac{1}{2-n} + \frac{1}{2-n'} - \frac{1}{2} \right]$$

$$\begin{aligned} \therefore \frac{1}{\mu^2} \left(\frac{dw}{dz} \right)^2 &= \left[\frac{1}{2-n} + \frac{1}{2-n'} - \frac{1}{2} \right]^2 \\ &= \left[\frac{1}{(2-n)^2} + \frac{1}{(2-n')^2} + \frac{1}{2^2} + \frac{2}{2(2-n)} - \frac{2}{2(2-n')} \right] \end{aligned}$$

from above we see that the function $f(z) = \frac{1}{\mu^2} \left(\frac{dw}{dz} \right)^2$ has poles at $z=2$ and $z=n'$ within C but the pole at $z=0$ is outside of C .

Hence $\text{Res}(z=0)$ is the sum of the co. of $\frac{1}{z}$ and hence ~~Res~~

$$\text{Res}(z=0) = \left[\frac{2}{2-n'} - \frac{2}{2-n} \right]_{z=0}$$

$$= 2 \left(\frac{1}{n'} + \frac{1}{n} \right)$$

Also $\text{Res}(z=n')$ = Sum of the co-ef. of $\frac{1}{z-n'}$ and hence

$$\begin{aligned} \text{Res}(z=n') &= \left[\frac{2}{2-n} - \frac{2}{2} \right]_{z=n'} \\ &= \frac{2}{n'-n} - \frac{2}{n'} \end{aligned}$$

So sum of the Residues within C is

$$= 2 \left(\frac{1}{n} + \frac{1}{n'} \right) + 2 \left(\frac{1}{n'-n} - \frac{1}{n'} \right)$$

$$= 2 \left(\frac{1}{n} + \frac{1}{n'-n} \right)$$

$$= \frac{2n'}{n(n'-n)}$$

$$= \frac{2a^2}{(a^2-n^2)n}$$

So, by Cauchy's Residue's theorem,

$$\int_C f(z) dz = \int_C \frac{1}{\mu^2} \left(\frac{dw}{dz} \right)^2 dz$$

$$= 2\pi i \times \text{sum of the residues}$$

$$= \frac{4\pi i a^2}{(a^2-n^2)n}$$

Also by Blasius theorem,

$$x - iy = \frac{ip}{2} \int_C \left(\frac{dw}{dz} \right)^2 dz$$

$$= \frac{ip}{2} \times \frac{4\pi a^2 i \mu^2}{(a^2-n^2)n}$$

Comparing with real and imaginary part,

$$x = - \frac{2\pi a^2 \mu^2 p}{(a^2-n^2)n} = \frac{2\pi a^2 \mu^2 p}{n(n-a^2)}$$

$$y = 0$$

So, resultant force acting on the circle
is $\sqrt{x^2 + y^2}$

$$= \frac{2Aa^2 \mu^2 p}{r(r-a^2)}$$

Above results shows that the force is
purely along OA.

So the disc will be ~~not~~ urged to move
along OA.

Also the cylinder is attracted ~~two~~ toward
the source and sketch of the stream
lines reveals that the pressure is great
on the opposite side of the disc than
that of the source.